

ON THE DYNKIN INDEX OF A PRINCIPAL  $\mathfrak{sl}_2$ -SUBALGEBRA

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## INTRODUCTION

The ground field  $\mathbb{k}$  is algebraically closed and of characteristic zero. Let  $\mathfrak{g}$  be a simple Lie algebra over  $\mathbb{k}$ . The goal of this note is to prove a closed formula for the Dynkin index of a principal  $\mathfrak{sl}_2$ -subalgebra of  $\mathfrak{g}$ , see Theorem 3.2. The key step in the proof uses the “strange formula” of Freudenthal–de Vries. As an application, we (1) compute the Dynkin index any simple  $\mathfrak{g}$ -module regarded as  $\mathfrak{sl}_2$ -module and (2) obtain an identity connecting the exponents of  $\mathfrak{g}$  and the dual Coxeter numbers of both  $\mathfrak{g}$  and  $\mathfrak{g}^\vee$ , see Section 4.

## 1. THE DYNKIN INDEX OF REPRESENTATIONS AND SUBALGEBRAS

Let  $\mathfrak{g}$  be a simple finite-dimensional Lie algebra of rank  $n$ . Let  $\mathfrak{t}$  be a Cartan subalgebra, and  $\Delta$  the set of roots of  $\mathfrak{t}$  in  $\mathfrak{g}$ . Choose a set of positive roots  $\Delta^+$  in  $\Delta$ . Let  $\Pi$  be the set of simple roots and  $\theta$  the highest root in  $\Delta^+$ . As usual,  $\rho = \frac{1}{2} \sum_{\gamma \in \Delta^+} \gamma$ . The  $\mathbb{Q}$ -span of all roots is a ( $\mathbb{Q}$ -)subspace of  $\mathfrak{t}^*$ , denoted  $\mathcal{E}$ . Choose a non-degenerate invariant symmetric bilinear form  $(\ , \ )_{\mathfrak{g}}$  on  $\mathfrak{g}$  as follows. The restriction of  $(\ , \ )_{\mathfrak{g}}$  to  $\mathfrak{t}$  is non-degenerate, hence it induces the isomorphism of  $\mathfrak{t}$  and  $\mathfrak{t}^*$  and a non-degenerate bilinear form on  $\mathfrak{t}^*$ . We require that  $(\theta, \theta)_{\mathfrak{g}} = 2$ , i.e.,  $(\beta, \beta)_{\mathfrak{g}} = 2$  of any long root  $\beta$  in  $\Delta$ .

**Definition 1** (E.B. Dynkin).

- (1) Let  $\mathfrak{s}$  be a simple subalgebra of  $\mathfrak{g}$ . The *Dynkin index* of  $\mathfrak{s}$  in  $\mathfrak{g}$  is defined by

$$\text{ind}(\mathfrak{s} \hookrightarrow \mathfrak{g}) = \frac{(x, x)_{\mathfrak{g}}}{(x, x)_{\mathfrak{s}}}, \quad x \in \mathfrak{s}.$$

- (2) If  $\nu : \mathfrak{g} \rightarrow \mathfrak{sl}(V)$  is a representation of  $\mathfrak{g}$ , then the *Dynkin index of the representation*, denoted  $\text{ind}_D(\mathfrak{g}, V)$  or  $\text{ind}_D(\mathfrak{g}, \nu)$ , is defined by

$$\text{ind}_D(\mathfrak{g}, V) = \text{ind}(\mathfrak{g} \hookrightarrow \mathfrak{sl}(V)).$$

It is not hard to verify that, for the simple Lie algebra  $\mathfrak{sl}(V)$ , the normalised bilinear form is given by  $(x, x)_{\mathfrak{sl}(V)} = \text{tr}(x^2)$ ,  $x \in \mathfrak{sl}(V)$ . Therefore, a more explicit expression for the Dynkin index of a representation  $\nu : \mathfrak{g} \rightarrow \mathfrak{sl}(V)$  is

$$(1.1) \quad \text{ind}_D(\mathfrak{g}, V) = \frac{\text{tr}(\nu(x)^2)}{(x, x)_{\mathfrak{g}}}.$$

Conversely, the index of a simple subalgebra can be expressed via indices of representations. Namely,

$$(1.2) \quad \text{ind}(\mathfrak{s} \hookrightarrow \mathfrak{g}) = \frac{\text{ind}_D(\mathfrak{s}, \mathfrak{g})}{\text{ind}_D(\mathfrak{g}, \text{ad}_{\mathfrak{g}})}.$$

The denominator in the right hand side represents the index of the adjoint representation of  $\mathfrak{g}$ , and the numerator represents the index of the  $\mathfrak{s}$ -module  $\mathfrak{g}$ .

The following properties easily follow from the definition:

**Multiplicativity:** If  $\mathfrak{h} \subset \mathfrak{s} \subset \mathfrak{g}$  are simple Lie algebras, then  $\text{ind}(\mathfrak{h} \subset \mathfrak{s}) \cdot \text{ind}(\mathfrak{s} \subset \mathfrak{g}) = \text{ind}(\mathfrak{h} \subset \mathfrak{g})$ .

**Additivity:**  $\text{ind}_D(\mathfrak{g}, V_1 \oplus V_2) = \text{ind}_D(\mathfrak{g}, V_1) + \text{ind}_D(\mathfrak{g}, V_2)$ . It is therefore sufficient to determine the indices for the irreducible representations.

**Theorem 1.1** (Dynkin, [2, Theorem 2.5]). *Let  $V_\lambda$  be a simple finite-dimensional  $\mathfrak{g}$ -module with highest weight  $\lambda$ . Then*

$$\text{ind}_D(\mathfrak{g}, V_\lambda) = \frac{\dim V_\lambda}{\dim \mathfrak{g}} (\lambda, \lambda + 2\rho)_{\mathfrak{g}}.$$

Although it is not obvious from the definition, the Dynkin index of a representation is an integer. This was proved by E.B. Dynkin [2, Theorem 2.2] using lengthy classification results. Later, he gave a better proof that is based on a topological interpretation of the index. A short algebraic proof is given in [5, Ch. I, §3.10].

**Example 1.2.**

1) Let  $R_d$  be the simple  $\mathfrak{sl}_2$ -module of dimension  $d + 1$ . Then  $\text{ind}_D(\mathfrak{sl}_2, R_d) = \binom{d+2}{3}$ .

2) Recall that  $\theta$  is the highest root in  $\Delta^+$ . By Theorem 1.1,

$$\text{ind}_D(\mathfrak{g}, \text{ad}) = (\theta, \theta + 2\rho)_{\mathfrak{g}} = (\theta, \theta)_{\mathfrak{g}} (1 + (\rho, \theta^\vee)_{\mathfrak{g}}) = 2(1 + (\rho, \theta^\vee)_{\mathfrak{g}}).$$

Note that the value  $(\rho, \theta^\vee)_{\mathfrak{g}}$  does not depend on the normalisation of the bilinear form. The integer  $1 + (\rho, \theta^\vee)$  is customary called the *dual Coxeter number* of  $\mathfrak{g}$ , and we denote it by  $h^*(\mathfrak{g})$ . Thus,  $\text{ind}_D(\mathfrak{g}, \text{ad}) = 2h^*(\mathfrak{g})$ . In the simply-laced case,  $h^*(\mathfrak{g}) = h(\mathfrak{g})$ —the usual Coxeter number. For the other simple Lie algebras, we have  $h^*(\mathbf{B}_n) = 2n-1$ ,  $h^*(\mathbf{C}_n) = n+1$ ,  $h^*(\mathbf{F}_4) = 9$ ,  $h^*(\mathbf{G}_2) = 4$ .

Andreev, Vinberg, and Elashvili applied the Dynkin index of representations to some invariant-theoretic problem [1]. To this end, they adjusted the index so that it does not depend on the choice of a bilinear form on  $\mathfrak{g}$ .

**Definition 2** (Andreev–Vinberg–Elashvili, 1967). Let  $\nu : \mathfrak{g} \rightarrow \mathfrak{sl}(V)$  be a finite-dimensional representation of a simple Lie algebra. Then

$$\text{ind}_{AVE}(\mathfrak{g}, V) := \frac{\text{ind}_D(\mathfrak{g}, V)}{\text{ind}_D(\mathfrak{g}, \text{ad})} = \frac{\text{tr}(\nu(x)^2)}{\text{tr}(\text{ad}_{\mathfrak{g}}(x)^2)}, \quad x \in \mathfrak{g}.$$

It follows that  $\text{ind}_{AVE}(\mathfrak{g}, \text{ad}_{\mathfrak{g}}) = 1$  and

$$\text{ind}_{AVE}(\mathfrak{g}, V_{\lambda}) = \frac{\dim V_{\lambda}}{\dim \mathfrak{g}} \cdot \frac{(\lambda, \lambda + 2\rho)_{\mathfrak{g}}}{(\theta, \theta + 2\rho)_{\mathfrak{g}}}.$$

## 2. THE “STRANGE FORMULA”

Let  $\mathcal{K}$  be the Killing form on  $\mathfrak{g}$ , i.e.,  $\mathcal{K}(x, x) = \text{tr}(\text{ad}_{\mathfrak{g}}(x)^2)$ ,  $x \in \mathfrak{g}$ . The induced bilinear form on  $\mathfrak{t}^*$  (and  $\mathcal{E}$ ) is denoted by  $\langle \cdot, \cdot \rangle$ . It is the so-called *canonical* bilinear form on  $\mathcal{E}$ . The canonical bilinear form is characterised by the following property:

$$(2.1) \quad \langle v, v \rangle = \sum_{\gamma \in \Delta} \langle v, \gamma \rangle \langle v, \gamma \rangle = 2 \sum_{\gamma > 0} \langle v, \gamma \rangle \langle v, \gamma \rangle \quad \text{for any } v \in \mathcal{E}.$$

The “strange formula” of Freudenthal–de Vries (see [3, 47.11]) is

$$\langle \rho, \rho \rangle = \frac{\dim \mathfrak{g}}{24}.$$

Using our normalisation of  $(\cdot, \cdot)_{\mathfrak{g}}$ , the “strange formula” reads

$$(2.2) \quad (\rho, \rho)_{\mathfrak{g}} = \frac{\dim \mathfrak{g}}{12} h^*(\mathfrak{g}).$$

Indeed, it is well known that  $\langle \theta, \theta \rangle = 1/h^*(\mathfrak{g})$  (see e.g. [6, Lemma 1.1]). Therefore, the transition factor between two forms  $\langle \cdot, \cdot \rangle$  and  $(\cdot, \cdot)_{\mathfrak{g}}$  (considered as forms on  $\mathcal{E}$ ) equals  $2h^*(\mathfrak{g})$ . Using the transition factor, we can also rewrite Eq. (2.1) in terms of  $(\cdot, \cdot)_{\mathfrak{g}}$ :

$$(2.3) \quad h^*(\mathfrak{g})(v, v)_{\mathfrak{g}} = \sum_{\gamma > 0} (v, \gamma)_{\mathfrak{g}}(v, \gamma)_{\mathfrak{g}}.$$

## 3. THE INDEX OF A PRINCIPAL $\mathfrak{sl}_2$ -SUBALGEBRA

If  $e \in \mathfrak{g}$  is nilpotent, then there exists a subalgebra  $\mathfrak{a} \subset \mathfrak{g}$  such that  $\mathfrak{a} \simeq \mathfrak{sl}_2$  and  $e \in \mathfrak{a}$  (Morozov, Jacobson). If  $e$  is a *principal* nilpotent element, then the corresponding  $\mathfrak{sl}_2$ -subalgebra is also called principal. (See [2, §9] and [4, Sect. 5] for properties of principal  $\mathfrak{sl}_2$ -subalgebras.) Let  $(\mathfrak{sl}_2)^{pr}$  be a principal  $\mathfrak{sl}_2$ -subalgebra of  $\mathfrak{g}$ . In this section, we obtain a uniform expression for  $\text{ind}((\mathfrak{sl}_2)^{pr} \hookrightarrow \mathfrak{g})$ .

Recall that  $\Delta$  has at most two root lengths. Let  $\theta_s$  denote the short dominant root in  $\Delta^+$ . (Hence  $\theta = \theta_s$  if and only if  $\Delta$  is simply-laced.) Set  $r = \|\theta\|^2/\|\theta_s\|^2 \in \{1, 2, 3\}$ . Along with  $\mathfrak{g}$ , we also consider the Langlands dual algebra  $\mathfrak{g}^{\vee}$ , which is determined by the dual root system  $\Delta^{\vee}$ . Since the Weyl groups of  $\mathfrak{g}$  and  $\mathfrak{g}^{\vee}$  are isomorphic, we have  $h(\mathfrak{g}) = h(\mathfrak{g}^{\vee})$ . However, the dual Coxeter numbers can be different (cf.  $\mathbf{B}_n$  and  $\mathbf{C}_n$ ).

The half-sum of positive roots for  $\mathfrak{g}^{\vee}$  is

$$\rho^{\vee} := \frac{1}{2} \sum_{\gamma > 0} \gamma^{\vee} = \sum_{\gamma > 0} \frac{\gamma}{(\gamma, \gamma)_{\mathfrak{g}}}.$$

It is well-known (and easily verified) that  $(\rho^\vee, \gamma)_{\mathfrak{g}} = \text{ht}(\gamma)$  for any  $\gamma \in \Delta^+$ . (This equality does not depend on the normalisation of a bilinear form.) It follows that  $h^*(\mathfrak{g}^\vee) = (\rho^\vee, \theta_s) = \text{ht}(\theta_s)$ .

**Proposition 3.1.** *For a simple Lie algebra  $\mathfrak{g}$  with the corresponding root system  $\Delta$ , we have*

$$(3.1) \quad \sum_{\gamma > 0} \text{ht}^2(\gamma) = \frac{\dim \mathfrak{g}}{12} h^*(\mathfrak{g}) h^*(\mathfrak{g}^\vee) r.$$

*Proof.* The equality in (3.1) is essentially equivalent to the “strange formula”.

Applying Eq. (2.3) to  $v = \rho^\vee$ , we obtain

$$(3.2) \quad h^*(\mathfrak{g})(\rho^\vee, \rho^\vee)_{\mathfrak{g}} = \sum_{\gamma > 0} (\rho^\vee, \gamma)_{\mathfrak{g}} (\rho^\vee, \gamma)_{\mathfrak{g}} = \sum_{\gamma > 0} \text{ht}^2(\gamma).$$

For  $\mathfrak{g}^\vee$ , the strange formula says that  $(\rho^\vee, \rho^\vee)_{\mathfrak{g}^\vee} = \frac{\dim \mathfrak{g}}{12} h^*(\mathfrak{g}^\vee)$ . Although the normalised bilinear forms  $(\ , \ )_{\mathfrak{g}}$  and  $(\ , \ )_{\mathfrak{g}^\vee}$  are proportional upon restriction to  $\mathcal{E}$ , they are not equal in general. Indeed, the square of the length of a long root in  $\Delta^\vee$  with respect to  $(\ , \ )_{\mathfrak{g}}$  equals  $2r$ . Hence the transition factor is  $r$  and

$$(3.3) \quad (\rho^\vee, \rho^\vee)_{\mathfrak{g}} = r(\rho^\vee, \rho^\vee)_{\mathfrak{g}^\vee} = \frac{\dim \mathfrak{g}}{12} h^*(\mathfrak{g}^\vee) r.$$

Then the assertion follows from (3.2) and (3.3).  $\square$

**Theorem 3.2.**  $\text{ind}((\mathfrak{sl}_2)^{pr} \hookrightarrow \mathfrak{g}) = \frac{\dim \mathfrak{g}}{6} h^*(\mathfrak{g}^\vee) r.$

*Proof.* Combining Eq. (1.2), Example 1.2(2), and Definition 2 yields the following formula for the index of a simple subalgebra  $\mathfrak{s}$  in  $\mathfrak{g}$ :

$$(3.4) \quad \text{ind}(\mathfrak{s} \hookrightarrow \mathfrak{g}) = \frac{h^*(\mathfrak{s})}{h^*(\mathfrak{g})} \cdot \text{ind}_{AVE}(\mathfrak{s}, \mathfrak{g}).$$

We use this formula with  $\mathfrak{s} = (\mathfrak{sl}_2)^{pr}$ . Let  $h$  be the semisimple element of a principal  $\mathfrak{sl}_2$ -triple. Without loss of generality, we may assume that  $h$  is dominant. Then  $\alpha(h) = 2$  for any  $\alpha \in \Pi$ . Put  $\tilde{h} = h/2$ . Then  $\gamma(\tilde{h}) = \text{ht}(\gamma)$  for any  $\gamma \in \Delta$  and  $\text{ad } \tilde{h}$  has the eigenvalues  $-1, 0, 1$  in  $(\mathfrak{sl}_2)^{pr}$ . Hence

$$\text{ind}_{AVE}((\mathfrak{sl}_2)^{pr}, \mathfrak{g}) = \frac{\text{tr}(\text{ad}_{\mathfrak{g}} \tilde{h})^2}{\text{tr}(\text{ad}_{\mathfrak{s}} \tilde{h})^2} = \frac{\sum_{\gamma \in \Delta} \text{ht}^2(\gamma)}{2} = \sum_{\gamma > 0} \text{ht}^2(\gamma).$$

Since  $h^*(\mathfrak{sl}_2) = 2$ , the theorem follows from Proposition 3.1 and Eq. (3.4).  $\square$

Below, we tabulate the values of index for all simple Lie algebras.

$\mathfrak{g}$	$\mathbf{A}_n$	$\mathbf{B}_n$	$\mathbf{C}_n$	$\mathbf{D}_n$	$\mathbf{E}_6$	$\mathbf{E}_7$	$\mathbf{E}_8$	$\mathbf{F}_4$	$\mathbf{G}_2$
$\text{ind}((\mathfrak{sl}_2)^{pr} \hookrightarrow \mathfrak{g})$	$\binom{n+2}{3}$	$\frac{n(n+1)(2n+1)}{3}$	$\binom{2n+1}{3}$	$\frac{(n-1)n(2n-1)}{3}$	156	399	1240	156	28

*Remark 3.3.* For the exceptional Lie algebras, Dynkin computed the indices of all  $\mathfrak{sl}_2$ -subalgebras, see [2, Tables 16–20].

Note that the index of a principal  $\mathfrak{sl}_2$  is preserved under the unfolding procedure  $\mathfrak{g} \rightsquigarrow \tilde{\mathfrak{g}}$  applied to multiply laced Dynkin diagram. Namely,  $\text{ind}((\mathfrak{sl}_2)^{pr} \hookrightarrow \mathfrak{g}) = \text{ind}((\mathfrak{sl}_2)^{pr} \hookrightarrow \tilde{\mathfrak{g}})$ , where the four pairs  $(\mathfrak{g}, \tilde{\mathfrak{g}})$  are:  $(\mathbf{C}_n, \mathbf{A}_{2n-1})$ ,  $(\mathbf{B}_n, \mathbf{D}_{n+1})$ ,  $(\mathbf{F}_4, \mathbf{E}_6)$ ,  $(\mathbf{G}_2, \mathbf{D}_4)$ . This is, of course, explained by the multiplicativity of the index of subalgebras and the fact that  $\text{ind}(\mathfrak{g} \hookrightarrow \tilde{\mathfrak{g}}) = 1$ .

*Remark 3.4.* Proposition 3.1 provides a uniform expression for  $\sum_{\gamma>0} \text{ht}^2(\gamma)$ . One might ask for a similar formula for  $\sum_{\gamma>0} \text{ht}(\gamma)$ . However, such a formula seems to only exist in the simply-laced case. Indeed, for any  $\mathfrak{g}$  we have  $2(\rho, \rho^\vee)_{\mathfrak{g}} = \sum_{\gamma>0} (\gamma, \rho^\vee)_{\mathfrak{g}} = \sum_{\gamma>0} \text{ht}(\gamma)$ . If  $\Delta$  is simply-laced, then  $\rho^\vee = 2\rho/(\theta, \theta)_{\mathfrak{g}} = \rho$ , and using the “strange formula” one obtains

$$\sum_{\gamma>0} \text{ht}(\gamma) = 2(\rho, \rho)_{\mathfrak{g}} = \frac{\dim \mathfrak{g}}{6} h(\mathfrak{g}) .$$

*Question.* Consider the function  $s \mapsto f(s) = \sum_{\gamma>0} \text{ht}^s(\gamma)$ . Are there some other values of  $s$  such that  $f(s)$  has a nice closed expression ?

#### 4. SOME APPLICATIONS

**(A)** Let  $\nu : \mathfrak{g} \rightarrow \mathfrak{sl}(V_\lambda)$  be an irreducible representation. Our first observation is that using Theorems 1.1 and 3.2 we can immediately compute the Dynkin index of  $V_\lambda$  as  $(\mathfrak{sl}_2)^{pr}$ -module:

$$\begin{aligned} \text{ind}_D((\mathfrak{sl}_2)^{pr}, V_\lambda) &= \text{ind}((\mathfrak{sl}_2)^{pr} \hookrightarrow \mathfrak{sl}(V_\lambda)) = \text{ind}((\mathfrak{sl}_2)^{pr} \hookrightarrow \mathfrak{g}) \cdot \text{ind}(\mathfrak{g} \hookrightarrow \mathfrak{sl}(V_\lambda)) = \\ \text{ind}((\mathfrak{sl}_2)^{pr} \hookrightarrow \mathfrak{g}) \cdot \text{ind}_D(\mathfrak{g}, V_\lambda) &= \frac{\dim \mathfrak{g}}{6} h^*(\mathfrak{g}^\vee) \cdot r \cdot \frac{\dim V_\lambda}{\dim \mathfrak{g}} (\lambda, \lambda + 2\rho)_{\mathfrak{g}} = \frac{\dim V_\lambda}{6} \cdot h^*(\mathfrak{g}^\vee) \cdot r \cdot (\lambda, \lambda + 2\rho)_{\mathfrak{g}} . \end{aligned}$$

Furthermore, we have

$$(4.1) \quad \text{ind}_D((\mathfrak{sl}_2)^{pr}, V_\lambda) = \text{ind}_D(\mathfrak{sl}_2, \text{ad}) \cdot \text{ind}_{AVE}((\mathfrak{sl}_2)^{pr}, V_\lambda) = 4 \cdot \text{ind}_{AVE}((\mathfrak{sl}_2)^{pr}, V_\lambda)$$

and

$$\text{ind}_{AVE}((\mathfrak{sl}_2)^{pr}, V_\lambda) = \frac{\text{tr}(\nu(\tilde{h})^2)}{\text{tr}((\text{ad } \tilde{h})^2)} = \frac{\sum_{\mu \vdash V_\lambda} \mu(\tilde{h})^2}{2} .$$

where notation  $\mu \vdash V_\lambda$  means that  $\mu$  is a weight of  $V_\lambda$ , and the sum runs over all weights according to their multiplicities. Since  $\mu(\tilde{h}) = (\mu, \rho^\vee)_{\mathfrak{g}}$ , we finally obtain

$$(4.2) \quad \sum_{\mu \vdash V_\lambda} (\mu, \rho^\vee)_{\mathfrak{g}}^2 = \frac{\dim V_\lambda}{12} \cdot h^*(\mathfrak{g}^\vee) \cdot r \cdot (\lambda, \lambda + 2\rho)_{\mathfrak{g}} .$$

This can be compared with the formula of Freudenthal–de Vries (see [3, 47.10.2]):

$$(4.3) \quad \sum_{\mu \vdash V_\lambda} \langle \mu, \rho \rangle^2 = \frac{\dim V_\lambda}{24} \langle \lambda, \lambda + 2\rho \rangle.$$

One can verify that Eq. (4.2) and (4.3) agree in the simply-laced case, where  $\rho$  is proportional to  $\rho^\vee$ .

**(B)** Let  $m_1, \dots, m_n$  be the exponents of  $\mathfrak{g}$ . Regarding  $\mathfrak{g}$  as  $(\mathfrak{sl}_2)^{pr}$ -module, one has  $\mathfrak{g} = \bigoplus_{i=1}^n R_{2m_i}$  [4, Cor. 8.7]. Then using Example 1.2(1), Eq. (3.4), (4.1), and the additivity of the index of representations, we obtain the identity

$$\begin{aligned} \frac{\dim \mathfrak{g}}{6} h^*(\mathfrak{g}^\vee) r = \text{ind}((\mathfrak{sl}_2)^{pr} \hookrightarrow \mathfrak{g}) &= \frac{h^*(\mathfrak{sl}_2)}{h^*(\mathfrak{g})} \sum_{i=1}^n \text{ind}_{AVE}(\mathfrak{sl}_2, R_{2m_i}) = \\ &= \frac{1}{2h^*(\mathfrak{g})} \sum_{i=1}^n \text{ind}_D(\mathfrak{sl}_2, R_{2m_i}) = \frac{1}{2h^*(\mathfrak{g})} \sum_{i=1}^n \binom{2m_i + 2}{3}. \end{aligned}$$

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